

# INTERMEDIATE-MASS HOT CORES AT $\sim 500$ AU: DISKS OR OUTFLOWS?

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## ABSTRACT

Observations with the Plateau de Bure Interferometer in the most extended configuration toward two intermediate-mass star-forming regions, IRAS 22198+6336 and AFGL 5142, reveal the presence of several complex organic molecules at  $\sim 500$  AU scales, confirming the presence of hot cores in both regions. The hot cores are not rich in CN-bearing molecules, as often seen in massive hot cores, and are mainly traced by CH<sub>3</sub>CH<sub>2</sub>OH, (CH<sub>2</sub>OH)<sub>2</sub>, CH<sub>3</sub>COCH<sub>3</sub>, and CH<sub>3</sub>OH, with additionally CH<sub>3</sub>CHO, CH<sub>3</sub>OD and HCOOD for IRAS 22198+6336, and C<sub>6</sub>H, and O<sup>13</sup>CS for AFGL 5142. The emission of complex molecules is resolved down to sizes of  $\sim 300$  and  $\sim 600$  AU, for IRAS 22198+6336 and AFGL 5142, respectively, and most likely is tracing protostellar disks rather than flattened envelopes or toroids as usually found. This is specially clear for the case of IRAS 22198+6336, where we detect a velocity gradient for all the mapped molecules perpendicular to the most chemically rich outflow of the region, yielding a dynamic mass  $\gtrsim 4 M_{\odot}$ . As for AFGL 5142, the hot core emission is resolved into two elongated cores separated  $\sim 1800$  AU. A detailed comparison of the complex molecule peaks to the new CO(2–1) data and H<sub>2</sub>O maser data from literature suggests that also for AFGL 5142 the complex molecules are mainly associated with disks, except for a faint and extended molecular emission found to the west, which is possibly produced in the interface between one of the outflows and the dense surrounding gas.

*Subject headings:* stars: formation — ISM: individual objects (IRAS 22198+6336, AFGL 5142) — ISM: lines and bands — radio continuum: ISM

## 1. INTRODUCTION

Hot molecular cores are compact ( $\leq 0.05$  pc) objects with high temperatures ( $\gtrsim 100$  K) and densities ( $n \gtrsim 10^6$  cm<sup>−3</sup>), which are characterized by a very rich chemistry of complex organic molecules (COMs, molecules with more than 6 atoms) and are typically associated with deeply embedded massive protostars (e. g., Cesaroni 2005). Such a rich chemistry is supposed to be triggered by the radiation from the nascent massive star which evaporates the complex molecules from the dust grain mantles and triggers additional gas-phase reactions (e. g., Millar et al. 1997). Thus, hot cores are supposed to be radiatively heated and originate in the innermost parts of the condensation where the massive star is being

formed. However, the radiative heating mechanism is questioned, as some observations suggest that the COM emission could be associated with shocks as well (Liu et al. 2002; Chandler et al. 2005; Goicoechea et al. 2006; Goddi et al. 2011a; Favre et al. 2011a; Zapata et al. 2011). This question cannot be easily answered because most of the hot cores known to date ( $\sim 100$ ) are associated with high-mass protostars, and are on average located at distances  $> 2$  kpc (studied typically at spatial scales  $\gtrsim 4000$  AU), while in the low-mass case (studied at scales 500–1000 AU, e. g., Bisschop et al. 2008), there are very few clear cases (the so-called hot corinos: e. g., Ceccarelli 2004; Bottinelli et al. 2007). Concerning the intermediate-mass regime, the only well-known cases are NGC 7129 (Fuente et al. 2005), IC 1396N (Fuente et al. 2009), and IRAS 22198+6336 (Sánchez-Monge et al. 2010), and still the presence of COMs is not clear in all the cases, questioning the true nature of these objects as hot cores. Thus, a detailed high spatial resolution ( $\sim 500$  AU) study of true hot cores (i. e., with complex organic chemistry) is lacking.

In this letter, we present new subarcsecond interferometric observations of COMs towards two intermediate-mass star-forming regions containing hot core candidates. The two regions are IRAS 22198+6336 (hereafter I22198), an intermediate-mass protostar of  $370 L_{\odot}$  located at 760 pc (Hirota et al. 2008), deeply embedded and driving a quadrupolar outflow with a thermal radiojet at its base (Sánchez-Monge et al. 2010), and AFGL 5142 (hereafter A5142), a protocluster of  $2300 L_{\odot}$  (Sánchez-Monge 2011) located at 1.8 kpc, driving at least 3 outflows and associated also with thermal free-free emission and with no infrared emission (Hunter et al. 1995). Submillimeter Array observations of CH<sub>3</sub>CN at  $\sim 2''$  angular resolution toward both regions indicate that the protostars are embedded in dense gas at around 100 K, but no clear evidence of COMs has been reported so far (Zhang et al. 2007;

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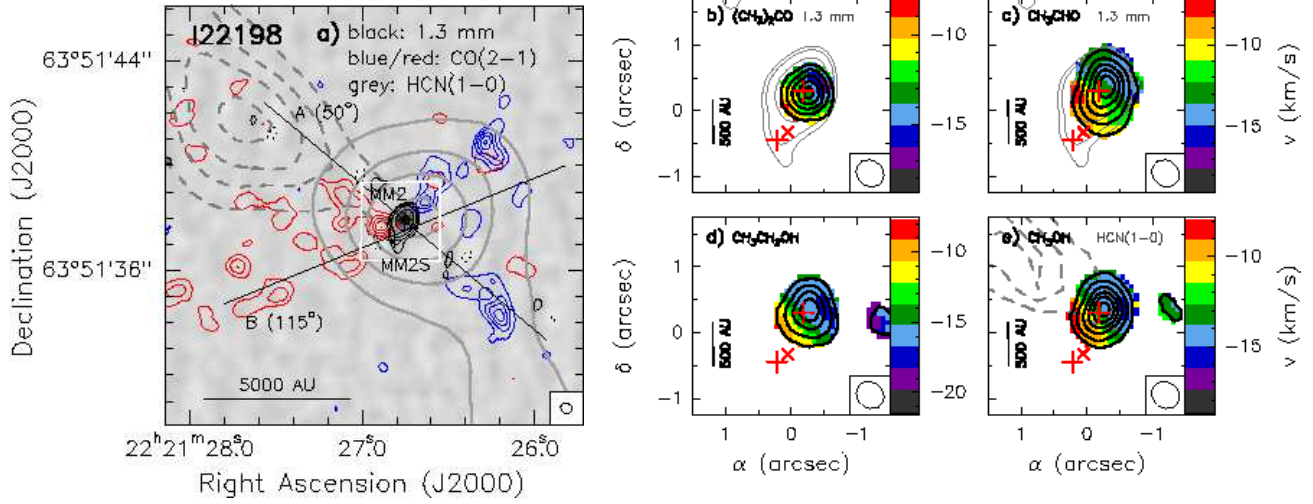


FIG. 1.— I22198 region. **a)** Greyscale and black contours are the 1.3 mm continuum emission (natural weight; beam:  $0.43'' \times 0.39''$ , P.A.= $27.1^\circ$ ; rms:  $2.0 \text{ mJy beam}^{-1}$ ). Contour levels are 4, 6, 10, 20, and 40 times  $2 \text{ mJy beam}^{-1}$ . Grey contours are the red(dashed) and blue(solid)-shifted emission of HCN (1–0) tracing outflow A (Sánchez-Monge 2011). Blue/red contours: CO (2–1) emission integrated from  $-31.0$  to  $-19.7/-5.3$  to  $7.9 \text{ km s}^{-1}$  (beam:  $0.55'' \times 0.50''$ , P.A.= $42.5^\circ$ , tapered to 400 m). Blue/red contour levels start at 20%/18% of the maximum,  $1754.97/1766.58 \text{ Jy beam}^{-1} \text{ m s}^{-1}$ , and increase in steps of 20%. **b–e)** Black thick contours/colorscales: zero/first-order moment of different COMs (at 230.176 (b), 230.316 (c), 230.231 (d), and 230.368 GHz (e); beam:  $0.44'' \times 0.39''$ , P.A.= $22.3^\circ$ ). Contour levels are 3, 10, 20, 30, 40, 50, and 60 times  $15 \text{ Jy beam}^{-1} \text{ m s}^{-1}$ . Black thin contours: **b–c)** 1.3 mm continuum emission. **e)** red-shifted HCN (1–0) emission (same as panel a). In all panels, red plus signs correspond to the millimeter continuum peaks, and the cross to the Spitzer-IRAC ( $3.6 \mu\text{m}$ ) source in the field (Sánchez-Monge, priv. commun.).

Sánchez-Monge et al. 2010). Here we show multiple detections of COMs down to angular scales of  $0.4''$ .

## 2. OBSERVATIONS

The IRAM Plateau de Bure Interferometer (PdBI) was used in its most extended configuration to observe the continuum and CO (2–1) emission at 230.538 GHz of I22198 and A5142, providing baselines in the range 136–760 m. A5142 was observed on 2010 January 10th (phase center:  $05:30:48.02 +33:47:54.5$ ), with system temperature of 170 K and averaged atmospheric precipitable water vapor around 2 mm. Bandpass calibration was carried out using B0851+202. To calibrate the phases we used B0552+398, and J0512+294, yielding a phase rms of  $20^\circ$ – $60^\circ$ . A subsequent track was carried out on 2010 January 18th to observe in track-sharing mode the I22198 region (phase center:  $22:21:26.78 +63:51:37.6$ ). The system temperature was  $\sim 150 \text{ K}$  and the precipitable water vapor was  $\sim 1 \text{ mm}$ . 3C454.3 was used for the bandpass calibration. Phases were calibrated with B2146+608, and B2037+511, and we obtained a phase rms of  $20^\circ$ – $40^\circ$ . The absolute flux density scale was determined from 3C273 and B2146+608, with an estimated uncertainty around 30%. The estimated uncertainty in absolute position is  $0.03''$ – $0.05''$ .

A correlator unit of 40 MHz bandwidth with 256 spectral channels was used to observe the CO (2–1) line in each polarization. Three additional units (per polarization) of 320 MHz bandwidth with 128 channels were used to observe the continuum across  $\sim 1 \text{ GHz}$ . The COMs were detected in the 320 MHz units, which provide a spectral resolution of 2.5 MHz or  $3.25 \text{ km s}^{-1}$ . Calibration and imaging were performed using the GILDAS software package<sup>13</sup>, following the standard procedures. The synthesized beams are  $0.43'' \times 0.39''$  at P.A.= $27.1^\circ$  for I22198, and  $0.50'' \times 0.30''$  at P.A.= $2.1^\circ$  for A5142. The rms of the final cleaned maps are given in the

captions of Figures 1 and 2.

## 3. RESULTS

### 3.1. Continuum and CO (2–1)

In Figures 1-a and 2-a we present the 1.3 mm continuum and CO (2–1) emission of I22198 and A5142 down to  $0.4''$ . For I22198 we detected one strong source, MM2 (following the nomenclature in Sánchez-Monge et al. 2010), extended in the southeast-northwest direction, with a faint extension at  $6\sigma$ , MM2-S,  $0.8''$  to the southeast. The overall extended emission of MM2 is perpendicular to the direction of outflow A (Sánchez-Monge et al. 2010). We fitted in the  $uv$ -plane an elliptical Gaussian to MM2 and obtained residual emission at the position of MM2-S, indicating an additional point source. The coordinates determined for MM2-S are (J2000):  $22:21:26.807$ ,  $63:51:37.14$ , and the flux density is  $18.2 \pm 0.8 \text{ mJy}$ , which corresponds to a mass of  $0.1$ – $0.6 M_\odot$ , assuming a dust temperature of 10–30 K, a gas-to-dust mass ratio of 100, and a dust mass opacity coefficient at 1.3 mm of  $0.899 \text{ cm}^2 \text{ g}^{-1}$  (agglomerated grains with thin ice mantles for densities  $\sim 10^6 \text{ cm}^{-3}$ , Ossenkopf & Henning 1994). The uncertainty in the masses is estimated to be a factor of 2. As for I22198-MM2, the deconvolved size is  $500 \times 300 \text{ AU}$  at P.A.= $-35^\circ$ , the peak intensity and flux density are  $91.8 \pm 0.9 \text{ mJy beam}^{-1}$ , and  $246 \pm 3 \text{ mJy}$ , and the mass is  $\sim 1 M_\odot$  (Table 1). Towards A5142 the millimeter emission is dominated by two partially extended and strong sources, MM1 and MM2, which are surrounded by five faint point-like sources (Palau et al., in prep.). The deconvolved sizes (from elliptical Gaussian fits in the  $uv$ -plane) are  $1200 \times 900 \text{ AU}$  at P.A.= $-86^\circ$  for MM1 and  $1000 \times 400 \text{ AU}$  at P.A.= $+18^\circ$  for MM2. The peak intensities and flux densities are  $38 \pm 3 \text{ mJy beam}^{-1}$ ,  $212 \pm 7 \text{ mJy}$  for MM1, and  $62 \pm 3 \text{ mJy beam}^{-1}$ ,  $151 \pm 4 \text{ mJy}$  for MM2, yielding masses of  $\sim 4 M_\odot$  (Table 1).

Regarding the CO (2–1) emission, we first caution that an important part of the emission is filtered out by the interferom-

<sup>13</sup> GILDAS: Grenoble Image and Line Data Analysis System, see <http://www.iram.fr/IRAMFR/GILDAS>.

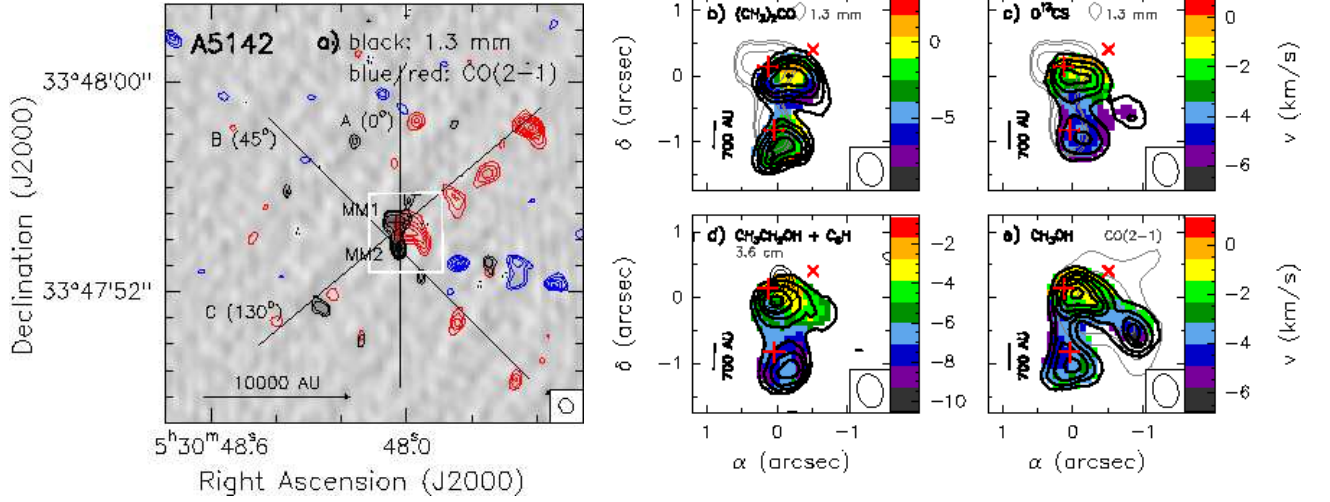


FIG. 2.— A5142 region. **a)** Greyscale and black contours are the 1.3 mm continuum emission (uniform weight; beam:  $0.50'' \times 0.30''$ , P.A. =  $2.1^\circ$ ; rms:  $2.8 \text{ mJy beam}^{-1}$ ). Contour levels are 4, 6, 10, 20, and 40 times  $2.4 \text{ mJy beam}^{-1}$ . Blue/red contours: CO(2–1) emission integrated from  $-19.7$  to  $-8.9/0.7$  to  $12.7 \text{ km s}^{-1}$  (beam:  $0.59'' \times 0.51''$ , P.A. =  $27.7^\circ$ , tapered to  $300 \text{ m}$ ). Blue/red contour levels start at 50%/37% of the maximum,  $1851.04/3175.57 \text{ Jy beam}^{-1} \text{ m s}^{-1}$ , and increase in steps of 10%. **b–e)** Black thick contours/colorscale: zero/first-order moment of different COMs (see Figure 1) (beam:  $0.52'' \times 0.36''$ , P.A. =  $12.0^\circ$ ). Contour levels are 4, 8, 12, 16, 20, and 24 times  $27 \text{ Jy beam}^{-1} \text{ m s}^{-1}$ . Black thin contours: **b–c)** 1.3 mm continuum emission; **d)** centimeter continuum (Zhang et al. 2007); **e)** CO(2–1) redshifted emission. Symbols as in Figure 1 (Spitzer source from Qiu et al. 2008).

eter and we are only sensitive to compact knots, even after tapering the data to a final beam of  $0.6''$ . In I22198 we detected chains of knots possibly tracing the cavity walls of outflows A and B (Sánchez-Monge et al. 2010). As for A5142, the CO(2–1) emission is again very clumpy but showing chains of knots which match well the known outflows of the region (e.g., Zhang et al. 2007).

### 3.2. Hot molecular core

In Figure 3-left we show the PdBI spectrum towards I22198-MM2, A5142-MM1 and A5142-MM2 for the entire observed bandwidth. Line identification was performed following the methodology described below. First, we searched for molecules with  $\geq 5$  transitions detectable within the observed frequency range and compared preliminary synthetic spectra (see below) for these molecules to the observed spectra. We found that, for the three sources,  $\text{CH}_3\text{CH}_2\text{OH}$  and  $(\text{CH}_2\text{OH})_2$  alone could account for about half of the detected transitions (Fig. 3-left). Second, we computed rotational diagrams (Fig. 4) with the  $\text{CH}_3\text{CH}_2\text{OH}$  transitions, allowing us to determine the gas temperature for each region (120 K for I22198, 210 K for A5142-MM1, and 140 K for A5142-MM2). Third, we adopted the derived temperatures, as well as the average linewidths of the involved transitions (given in Fig. 3), to compute synthetic spectra for both  $\text{CH}_3\text{CH}_2\text{OH}$  and  $(\text{CH}_2\text{OH})_2$ , summed them, and subtract the sum to the observed spectrum. Finally, the definitive line identification of molecules different from  $\text{CH}_3\text{CH}_2\text{OH}$  and  $(\text{CH}_2\text{OH})_2$  was performed in the residual spectrum (Fig. 3-right) for lines with flux above  $4\sigma$  ( $\sigma = 4.5 \text{ mJy}$  for I22198;  $\sigma = 7.5 \text{ mJy}$  for A5142). The systemic velocity used for I22198 was derived from the strongest isolated lines and was found to be  $-12.3 \text{ kms}$ . As for A5142, we adopted the systemic velocities derived by Zhang et al. (2007):  $-1.0 \text{ km s}^{-1}$  for A5142-MM1, and  $-3.4 \text{ km s}^{-1}$  for A5142-MM2. The synthetic spectra were computed using the estimates of temperature and linewidth given above, and assuming local thermodynamic equilibrium, optically thin emission, and the molecular data from the Jet

Propulsion Laboratory (Pickett et al. 1998) or the Cologne Database for Molecular Spectroscopy catalogs (Müller et al. 2005), except for  $\text{CH}_3\text{OD}$  (Anderson et al. 1988). The final identification is presented in Fig 3-right and the column densities used to build the synthetic spectra are listed in Table 1.

In addition to  $\text{CH}_3\text{CH}_2\text{OH}$  and  $(\text{CH}_2\text{OH})_2$ , the strongest lines found in the spectra of the three sources are from  $\text{CH}_3\text{OH}$  and  $\text{CH}_3\text{COCH}_3$ , and for none of the sources CN-bearing species such as  $\text{CH}_2\text{CHCN}$ ,  $\text{CH}_3\text{CH}_2\text{CN}$  were required to fit the spectra. Molecules detected only in I22198 are  $\text{CH}_3\text{CHO}$ ,  $\text{CH}_3\text{OCHO}$ ,  $\text{CH}_3\text{OD}$  and  $\text{HCOOD}$ . On the other hand, for A5142-MM1, no deuterated species were detected,  $\text{O}^{13}\text{CS}$  was found to dominate the line at  $230.317 \text{ GHz}$  (with no need of  $\text{CH}_3\text{CHO}$ ), and we identified two transitions of  $\text{C}_6\text{H}$ , with energies of the upper state of 483 K. The spectrum observed in A5142-MM2 is essentially the same as that of A5142-MM1, with smaller fluxes.

Among the strongest detected transitions we have chosen four to be representative of different excitation conditions, and computed the zero-order (integrated intensity) and first-order (velocity) maps (Figures 1-b–e and 2-b–e). From the figure it is seen that in I22198 the COM emission is restricted to MM2 without extending to MM2-S, and that this emission is elongated for almost all the molecules in the southeast-northwest direction (the only unresolved emission is that from  $\text{CH}_3\text{COCH}_3$ , as in Orion-KL, e.g., Friedel et al. 2005). The deconvolved size and P.A. of the average emission of the four transitions shown in Fig. 1-b–e is listed in Table 1. It is interesting to note that the first-order moment map of the three resolved COM transitions of I22198 shows a velocity gradient in the direction perpendicular to outflow A (Fig. 1e). Concerning A5142, COM emission is found in both millimeter continuum sources MM1 and MM2, with the first-order moment showing that there is a shift in velocities between the two sources of  $\sim 3 \text{ km s}^{-1}$ , as found by Zhang et al. (2007), and with hints of an elongation in the east-west direction for MM1, and in the southeast-northwest direction for MM2 (Table 1). What is more, the emission from MM1 reveals an ex-



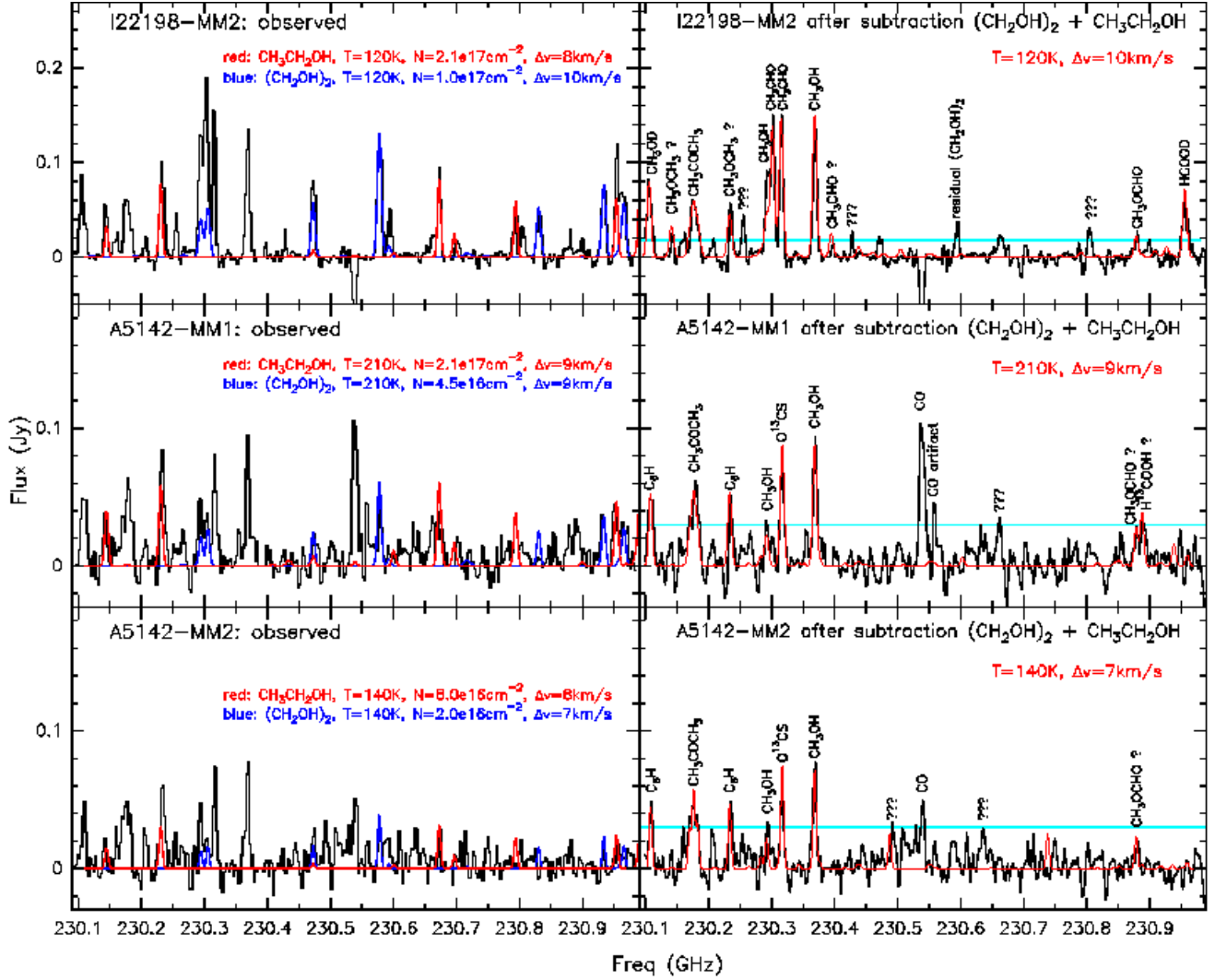


FIG. 3.— *Left*: Spectrum for I22198 (top), A5142-MM1 (middle), and A5142-MM2 (bottom) over the PdBI continuum bands, with the computed synthetic spectra overlaid for  $\text{CH}_3\text{CH}_2\text{OH}$  and  $(\text{CH}_2\text{OH})_2$  (see text). The spectrum is extracted by integrating the intensity over an area of  $0.41''$ ,  $0.52''$ ,  $0.43''$  (FWHM), yielding conversion factors of  $136.6$ ,  $84.9$ , and  $124.2 \text{ K Jy}^{-1}$  for I22198, A5142-MM1, A5142-MM2, respectively. *Right*: Residual spectra after subtracting the contribution from  $\text{CH}_3\text{CH}_2\text{OH}$  and  $(\text{CH}_2\text{OH})_2$ , with the identified lines marked at the top of each transition. The assumed temperatures and linewidths used to build the synthetic spectra are given in each panel, and the column densities used are listed in Table 1. The horizontal line indicates the  $4\sigma$  threshold used for line identification.

tension to the southwest, apparent mainly in  $\text{CH}_3\text{OH}$ , which is following the CO redshifted emission shown in Fig. 2-e.

#### 4. DISCUSSION AND CONCLUSIONS

##### 4.1. Hot corinos or massive hot cores?

The set of detected molecules in the intermediate-mass hot cores of I22198 and A5142 does not include CN-bearing species. We inspected the observed frequency range and found that at least two transitions of  $\text{CH}_2\text{CHCN}$  should have been detected (at  $230.488$  and  $230.739 \text{ GHz}$ ) with the same intensity, from synthetic spectra at rotational temperature in the range  $100$ – $600 \text{ K}$  and column densities  $\sim 10^{15} \text{ cm}^{-2}$  (Table 1). Thus, the set of detected molecules is similar to the sets of the hot corino IRAS 16293–2422, where  $\text{O}^{13}\text{CS}$ ,  $\text{CH}_3\text{CHO}$ , and  $\text{CH}_3\text{OCHO}$  are also detected (Bottinelli et al. 2007; Caux et al. 2011). However, IRAS 16293–2422 is dominated by simple O-rich and HCO-rich species like  $\text{H}_2\text{CO}$ ,  $\text{SO}_2$ ,  $\text{CH}_3\text{OH}$ ,  $\text{CH}_3\text{CHO}$ , and  $\text{CH}_3\text{OCHO}$ , while in I22198 and A5142 we detected more  $\text{CH}_2/\text{CH}_3$ -rich molecules, such as  $\text{CH}_3\text{CH}_2\text{OH}$ ,

$(\text{CH}_2\text{OH})_2$ , and  $\text{CH}_3\text{COCH}_3$ . On the other hand, massive hot cores, such as Orion-KL (Caselli et al. 1993; Blake et al. 1996; Friedel & Snyder 2008), G29.96–0.02 (Beuther et al. 2007), or G34.26+0.15 (Mookerjee et al. 2007) show strong emission in  $\text{CH}_2\text{CHCN}$ , and  $\text{CH}_3\text{CH}_2\text{CN}$ , as well as in  $\text{CH}_3\text{COCH}_3$  and  $\text{CH}_2/\text{CH}_3$ -rich molecules (e.g., Öberg et al. 2011). Thus, the intermediate-mass hot cores in I22198 and A5142 show a chemistry with properties of both low-mass hot corinos (lack of CN-bearing molecules) and massive hot cores ( $\text{CH}_2/\text{CH}_3$ -rich complex molecules).

##### 4.2. Hot core emission at $\sim 500 \text{ AU}$ : disks or outflows?

The sizes (FWHM) of the COM emission range from  $150$ – $300 \text{ AU}$  for I22198, and  $300$ – $1000 \text{ AU}$  for A5142 (Table 1), and are, after Orion-KL ( $200$ – $1000 \text{ AU}$ : Guelin et al. 2008, Favre et al. 2011a, 2011b; Zapata et al. 2011), among the smallest in intermediate/high-mass star-forming regions ever measured ( $> 2000 \text{ AU}$ : e.g., Beuther et al. 2007; Liu et al. 2010; Zapata et al. 2010; Beltrán et al. 2011). Thus, the

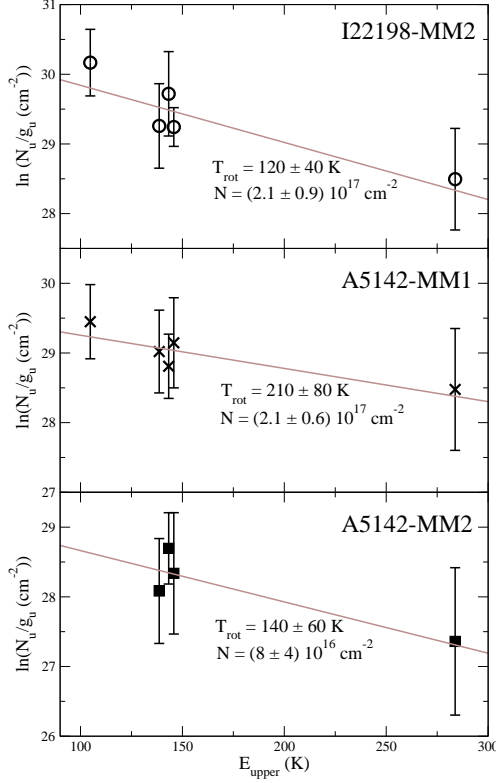


FIG. 4.— Rotational diagram for  $\text{CH}_3\text{CH}_2\text{OH}$  for I22198-MM2 (top), A5142-MM1 (middle), and A5142-MM2 (bottom). For those  $\text{CH}_3\text{CH}_2\text{OH}$  transitions which are significantly contaminated ( $\sim 50\%$ ) by emission from another molecule, half the line area is used (blended lines: 230.954 GHz (with energy of the upper state,  $E_u$ , of 146 K) for I22198, and 230.231 GHz ( $E_u = 143$  K) for A5142-MM1 and MM2.

present observations are well suited to consider whether the emission is due to internal heating (disks) and/or shocks (from outflows).

For the case of I22198, the continuum and COM emission peaks are coincident within  $< 0.1''$  (Table 1), and COM emission is not present near MM2-S, which in turn is coincident with an IRAC source at  $3.6 \mu\text{m}$  and  $4.5 \mu\text{m}$ . This is suggestive of MM2 being less evolved than MM2-S. On the other hand, the two outflows present in the region show different properties, with HCN and SiO emission detected only along the direction of outflow A (Fig. 1a, Sánchez-Monge 2011), which supports the interpretation of two different outflows against one single wide-angle outflow. Thus, a plausible interpretation is that MM2 is the driving source of outflow A. If this is the case, the COM emission, with a size of 300 AU, would be tracing a rotating disk, similar to the disk toward the NGC 1333-IRAS2A hot corino, of 200 AU of size (Joergensen et al. 2005). The velocity gradient seen in the COM emission of MM2 yields a central protostellar mass of  $\gtrsim 4.2 M_\odot$  (following Sánchez-Monge et al. 2010; lower limit from assuming an inclination angle with respect to the plane of the sky of  $90^\circ$ ), supporting the intermediate-mass nature of the object. From a theoretical point of view, models of disks in massive protostars predict disk sizes of hundreds of AU, and linewidths around  $10 \text{ km s}^{-1}$  (e. g., Krumholz et al. 2007), both parameters coincident with our measurements. Thus, although the COM emission could be contaminated by shocks from the outflows, our observations in I22198 are consistent with the COM emission arising from a disk.

TABLE 1  
PROPERTIES OF INTERMEDIATE-MASS HOT CORES AT  $\sim 500$  AU

	I22198-MM2	A5142-MM1	A5142-MM2
distance (pc)	760		1800
$L_{\text{bol}} (L_\odot)$	370		2300
<i>continuum + outflow information</i>			
clustering (mm)	no		yes
$L_{\text{cm}}$ (mJy kpc $^2$ )	0.34	2.1	$< 0.9$
outflow(s) PA ( $^\circ$ )	50	-40, 0	45
$M_{\text{env}} (M_\odot)^a$	0.5–2	2–9	2–6
$N_{\text{H}_2} (\text{cm}^{-2})^a$	$2 \times 10^{25}$	$1 \times 10^{25}$	$2 \times 10^{25}$
density ( $\text{cm}^{-3})^a$	$4 \times 10^9$	$1 \times 10^9$	$3 \times 10^9$
1.3 mm cont., RA	22:21:26.753	05:30:48.031	05:30:48.024
1.3 mm cont., DEC	63:51:37.90	33:47:54.64	33:47:53.67
continuum size (AU) $^b$	$520 \times 300$	$1200 \times 900$	$1000 \times 400$
continuum PA ( $^\circ$ ) $^b$	-35	-86	+18
<i>molecular hot core (HC) information<math>^c</math></i>			
$T_{\text{rot,CH}_3\text{CH}_2\text{OH}}$ (K)	$\sim 120$	$\sim 210$	$\sim 140$
$N_{\text{CH}_3\text{CH}_2\text{OH}} (\text{cm}^{-2})$	$2.1 \times 10^{17}$	$2.1 \times 10^{17}$	$8.0 \times 10^{16}$
$N_{(\text{CH}_2\text{OH})_2} (\text{cm}^{-2})$	$1.0 \times 10^{17}$	$4.5 \times 10^{16}$	$2.0 \times 10^{16}$
$N_{\text{CH}_3\text{COCH}_3} (\text{cm}^{-2})$	$5.0 \times 10^{16}$	$1.0 \times 10^{17}$	$4.0 \times 10^{16}$
$N_{\text{CH}_3\text{OH}} (\text{cm}^{-2})$	$2.3 \times 10^{19}$	$2.3 \times 10^{18}$	$4.0 \times 10^{18}$
$N_{\text{CH}_3\text{CHO}} (\text{cm}^{-2})$	$6.0 \times 10^{16}$	$< 1 \times 10^{16}$	$< 7 \times 10^{15}$
$N_{\text{HCOOD}} (\text{cm}^{-2})$	$3.0 \times 10^{16}$	$< 2 \times 10^{16}$	$< 1 \times 10^{16}$
$N_{\text{CH}_3\text{OCHO}} (\text{cm}^{-2})$	$2.5 \times 10^{17}$	$< 2 \times 10^{17}$	$< 2 \times 10^{17}$
$N_{\text{C}_6\text{H}} (\text{cm}^{-2})$	$< 6 \times 10^{15}$	$3.0 \times 10^{15}$	$5.0 \times 10^{15}$
$N_{\text{O}^{13}\text{CS}} (\text{cm}^{-2})$	$< 5 \times 10^{17}$	$1.5 \times 10^{16}$	$1.0 \times 10^{16}$
$N_{\text{CH}_2\text{CHCN}} (\text{cm}^{-2})$	$< 2 \times 10^{15}$	$< 3 \times 10^{15}$	$< 3 \times 10^{15}$
HC offset $^d$	-0'08,+0'01	-0'27,-0'20	-0'12,-0'21
HC size (AU) $^b$	$280 \times 160$	$1000 \times < 380$	$820 \times 340$
HC PA ( $^\circ$ ) $^b$	+155	+89	+120

$^a$  Masses computed assuming a dust temperature in the range 30–100 K, and the dust opacity law given in the main text. The  $\text{H}_2$  column density and density, estimated assuming spherical symmetry and uniform density, are calculated adopting a dust temperature of 50 K, and the mass and size from the continuum emission.

$^b$  Deconvolved size and P. A. For the case of the COMs, the (average) size is obtained from Gaussian fits to the zero-order moment map integrated in the range 230.1–230.4 GHz.

$^c$  Column densities are calculated assuming the rotational temperature and linewidth given in Fig. 4, and are corrected for filling factor (source sizes from Fig. 3).

$^d$  RA and DEC offset with respect to continuum peak position.

Concerning A5142, the COM emission is offset from the continuum emission by  $\sim 0.3''$  (or  $\sim 500$  AU, Table 1) for both MM1 and MM2. Possible explanations are: i) COM emission is optically thick; ii) COM emission is strongly affected by shocks and the passage of outflow(s); iii) the continuum and COM peaks are tracing different objects. We discard the first possibility because the peak of COMs for which a rotational diagram was performed is offset from the continuum peak and the rotational diagram does not show a clear sign of high opacity (Fig. 4). The second possibility of COM emission produced by shocks/outflows is questioned by the fact that the peak of COMs near A5142-MM1 is coincident with the peak of a centimeter source (Fig. 2-d), and the centimeter source, most likely tracing a thermal radiojet, falls exactly at the center of symmetry of outflow C, as traced by  $\text{H}_2\text{O}$  masers (see Figure 3 in Goddi & Moscadelli 2006). This suggests that the COM emission near MM1 could be associated with a possible disk (perpendicular to the centimeter source elongation) from the driving source of outflow C, which is consistent with recent  $\text{CH}_3\text{OH}$  and  $\text{H}_2\text{O}$  maser observations (Goddi et al. 2011b). In addition, the CO (2–1) emission of outflow

B seems to have its center of symmetry in MM2 (Fig. 2-a), as already suggested by Goddi & Moscadelli (2006), and interestingly the COM emission near MM2 is elongated in the direction perpendicular to outflow B (Fig. 2-b,d; Table 1), suggesting again the association with a disk.

To further assess the nature of the COM emission in A5142-MM1 and A5142-MM2 we have computed the ratio of a HCO-rich species with respect to a CH<sub>2</sub>/3-rich species, expected to be small for processed (disk) material (e. g., Öberg et al. 2011). The CH<sub>3</sub>CHO/CH<sub>3</sub>CH<sub>2</sub>OH ratio is < 0.05 for A5142-MM1, and < 0.09 for A5142-MM2, about one order of magnitude smaller than those derived in shocks associated with outflows, and similar to the values derived in hot corinos and hot cores (e. g., Arce et al. 2008; Öberg et al. 2011). This gives further support to the interpretation of COM emission in A5142-MM1 and MM2 being associated with disks, although obviously some (minor) contribution from shocks/outflows could still be present. Finally, the elongation of A5142-MM1 to the southwest could be arising in shocks/outflows, given its strong similarity to the CO (2–1) redshifted lobe (Fig. 2-

e), and the fact that its emission is detected in O<sup>13</sup>CS and CH<sub>3</sub>OH, which are molecules typically found in chemically rich outflows (e. g., Bachiller & Pérez-Gutiérrez 1997; Leurini et al. 2011).

In conclusion, the high angular resolution PdBI observations have revealed that the hot cores associated with the intermediate-mass protostars I22198 and A5142 are most likely tracing disks, with some contribution from the shocked gas of the outflows.

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